

Population Balance Modeling for Integrating Hydrodynamics with Electrostatics in Crude Oil Electrocoalescers

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ABSTRACT

A particular population balance model (PBM) which consists of hydrodynamic and electrostatic parts is developed for electrocoalescence of distributed water droplet in the continuous oil phase. The approach is the modification of a recognized PBM by adding the electrostatic effects on the overall coalescence rate including the number of times (frequency) occurring collision and the efficiency of coalescence. Moreover, the modified model has been being implemented in a CFD-PBM problem for a pilot plant electrocoalescer to predict the profile of water phase and size distribution of droplets. The results recognize the effect of local electric field intensity and local water content on electrocoalescence rate. Furthermore, the results demonstrate that separation for the very small droplets ($<4\mu\text{m}$) is minor, for the medium sizes ($8\text{--}32\mu\text{m}$) is more considerable, and for larger droplets ($>64\mu\text{m}$) occurs completely. Ultimately, by making a comparison between the simulation results and the pilot data, the EHD PB model is validated.

Keywords: Electrocoalescence, w/o Emulsion, Droplet, Electrohydrodynamic (EHD), Population Balance (PB).

INTRODUCTION

The flow of the produced crude oil and brine through reservoirs and pressure relief valves results in a water-in-oil emulsion. The conventional method which has been used to separate water from crude oil is mainly based on the gravity in which the density difference gets the water droplets separated as the heavier phase. According to the Stokes equation, the rate of settling process is proportional to the square of the droplet size and is controlled by the smallest droplets [1]. The most effective technique for speeding up the separation

process is to apply an electric field to force small water droplets to coalesce [2]. The electric field gives rise to attractive forces among the droplets and increases the probability of the coalescence [3,4]. Conventionally, this electrostatic separation method is operated in horizontal vessels well-known as electrocoalescers.

Water-in-oil emulsion, which consists of a wide range size of brine droplets, is dispersed in crude oil and comes into the electrocoalescer. Usually, the inlet feed is distributed uniformly through the orifices of the horizontal pipe which are placed at a

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designed elevation from the bottom of the vessel. The emulsion enters into the treater as laminar flow and moves upward to the strong electric field, where small water droplets are merged into larger ones. The effects of the electrostatic field can be explained by body forces acting on water droplets [5,6].

Crude oil electrocoalescer is a multiphase flow system which consists of water droplets dispersed in a continuous oil phase. Also, the main concern of the subject is to model droplet-scale phenomena like droplet-droplet collision and coalescence either forced by fluid flow or caused by electrostatic forces. Because of the complex behavior of water droplets from the simultaneous effects of electrostatic and hydrodynamic phenomena, there is not a comprehensive model for the electrocoalescers [7,8]. Therefore, a proper design of crude oil electrocoalescers requires pilot tests which are time and money consuming phase through the design stage. Furthermore, scale-up of an electrocoalescer from pilot scale to full-scale is not precise and requests skills and know-how.

Population Balance Modeling (PBM) and Computational Fluid Dynamics (CFD) are techniques that help us improve process equipment performance. In addition, CFD can determine fluid dynamics in process systems. Moreover, CFD provides the completely resolved flow field in equipment. PBM can identify coalescence, breakage, and size distribution of droplets while it considers the droplet-scale phenomena in the electrocoalescer. An increase in scale from a Lab-scale to a pilot scale can be effectively acquire experiments on binary droplet electrocoalescence. Scale-up from a pilot plant to a full-scale plant can be attained using PBM. Furthermore, PBM coupled

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with CFD can lead to high precision modeling of industrial electrocoalescers [9].

The interactions among droplets such as coalescence cause a distributed population of the droplets. Droplet-droplet coalescence may occur after the collision of two droplets either by electrostatic attractive forces or through flows of the continuous phase. The success coalescence among the collided droplets is completed if the trapped film is drained out. Moreover, most coalescence models are based on this film drainage theory [10-12]. On the other hand, some models were developed for the breakage rate of droplets by the same researchers [13]. Generally, droplet breakage arises from the effect of turbulent continuous phase on a droplet, if the turbulent kinetic energy overcomes on the surface energy. Furthermore, the breakage is expected to occur if the electric field exceeds the critical strength in electrocoalescers [14-16].

Both PBM and CFD have limitations when they are implemented individually. To run PBM, some information is needed from fluid dynamics of the system such as droplet velocity and turbulence which must be estimated by simplified correlations, unavoidably. In CFD simulations of a multiphase system, generally an equivalent diameter is used instead of the different sizes of the droplets. If the distribution of the droplet size is broad, this method is not accurate because coalescence or breakage is often not assessed [17]. In practice, a wide droplet size distribution exists in the conventional electrostatic dehydrators of the petroleum industry. Also, smallest droplets can be found near the oil outlet, while a strong coalescence rate under the effect of the electric field leads to very bigger droplets in the lower

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regions near the oil-water interface. These differences in the local size distribution of water droplets affect the hydrodynamics and explain changes in velocities and volume percentage of the water phase. Assuming a constant Sauter mean diameter cannot indicate these influences, and the resultant sophisticated behavior has been easier and less complicated in CFD simulations without PBM. So coupling CFD with PBM can share the advantages of two approaches [17].

A coalescence efficiency model which includes the attractive van der Waals and repulsive electrostatic force has been developed according to the DLVO theory into the PBE framework by Kamp and Kraume [18]. It is attempted by them to extend hydrodynamic models of coalescence efficiency so that the coalescence inhibition due to natural electrostatic interactions can be contributed. For this purpose, the model of Coulaloglou and Tavlaridec [10] with the DLVO models of Israelachvili [19] and Derjaguin et al. [20] have been coupled by them. In addition, considering independency of the two coalescence processes (hydrodynamic and electrostatic), the total coalescence efficiency as the product of these two efficiencies has been described by Derjaguin et al. The model of Kamp and Kraume, [18] which is defined as the Debye length, is limited to natural electrostatic interactions which is effective only at very close distances between two drops. Therefore, it is not applicable for electrocoalescence in which an external strong electrostatic field induces intense dipolar attraction force between two far drops. A combined lossy capacitor population balance model (LCPBM) has been developed by Barega et al [21] to predict the effect of a square wave frequency on electric-field coalescence. An electrical-circuit model has been

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coupled with the hydrodynamic equation by them to calculate the velocities of drops, which in turn (it has been) used in PBM to determine the time evolution of the drop size. However, it is expressed by them that their model is not able to consider drop-drop coalescence, and it only presents the overall final size distribution.

An AC electrostatic desalter in the steady state condition has been modeled by Meidanshahi et al [22] using a bivariate population balance equation. However, it seems that their model can simplify the size distribution of drops by assuming three types of emulsions as tight, medium, and loose; in addition, it only presents water content at the outlet. Therefore, their model is unable to predict time-dependent size distribution profiles of drops through the vessel. A mathematical model has been developed by Aryafard et al [23,24] based on population balance method for an industrial desalting plant at steady state condition. Furthermore, an empirical correlation developed by Manga and Stone [25] has been used in their model to consider the effect of the electric field, rather than a fundamental dynamic theorem. Therefore, their model is unable to predict time-dependent size distribution profiles of the drops. The existing PB models only consider hydrodynamic aspects. In other words, they cannot consider hydrodynamic aspects and electrostatic issues simultaneously. For the liquid-liquid subject in the region of electrocoalescence, no combined CFD-PBM model has been investigated yet. Moreover, only an electrostatic PBM has been presented by Atten [26]. Nevertheless, some explicit and implicit assumptions such as the approximation of homogeneous and isotropic spatial distribution of droplets, as well as the simplified hydrodynamic

aspect, limits the model of Atten to generalize directly for practical applications. Therefore, developing an electrohydrodynamic model with the consideration of the important microscale phenomena in electrocoalescers assists in predicting the internal situations of the phases and size distribution of droplets. Moreover, the resulted data achieved from CFD simulation by using the developed model will be so useful to design and improve the internal electric field and hydraulic of the fluid flow.

EXPERIMENTAL PROCEDURE

Population Balance Modeling

A modified droplet-droplet coalescence kernel is developed to describe the EHD phenomena in a w/o emulsion within an electrocoalescer. Because of very laminar flow at the inlet, as the emulsion flows in the strong electric field, probability of droplet breakup is negligible in comparison with droplet-droplet coalescence. Furthermore, the coalescence is less achievable by individual fluid flow impacts, but it is more possible under the effect of an electrostatic field [8].

Nevertheless, no practiced PB model has been developed for electrocoalescence. Therefore, a special PBM of coupled hydrodynamics and electrostatics has to be modified to consider both features simultaneously. The practical approach of this study is to modify an appropriate recognized hydrodynamic PBM in such a way to add the electrostatic effects on overall coalescence rate. Generally, the population balance equation is a continuity term from which could be obtained as a balance for droplets in an element volume. For a number density function $n(x, r, t)$ with the coordinates which are external (x) and internal (r),

the population balance equation could be as follows [27]:

$$\frac{\partial n}{\partial t} + \nabla_x \cdot \dot{X}n + \nabla_r \cdot \dot{R}n = p \quad (1)$$

where p in the above equation (Equation 1) is the net rate of production of droplets due to birth and death processes, and X and R are the velocities for external and internal coordinates respectively. In the present case, the continuity equation (Equation 2) for the number density function $n(x, V, t)$, which denotes the spatial position x of a droplet as an external coordinate and the droplet volume V as an internal coordinate, is given by:

$$\frac{\partial}{\partial t}[n(V, t)] + \nabla \cdot [\bar{u} n(V, t)] = S(V, t) \quad (2)$$

where $S(V, t)$ is the source term due to coalescence and breakup of droplets of volume V . Furthermore, it can be written as:

$$S(V, t) = B^c(V, t) - D^c(V, t) + B^b(V, t) - D^b(V, t) \quad (3)$$

B^c , D^c , B^b , and D^b are the birth and death rates of droplets of volume V due to coalescence and breakup. Generally, the birth and death rates, B^c and D^c of droplets of volume V , due to coalescence are given by the following equations:

$$B^c(V, t) = \frac{1}{2} \int_0^V c(V - V', V') n(V - V', t) n(V', t) dV' \quad (4)$$

$$D^c(V, t) = \int_0^\infty c(V, V') n(V, t) n(V', t) dV' \quad (5)$$

where $c(V, V')$ is the coalescence rate of droplets of volume V with droplets of volume V' .

From the hydrodynamic point of view, the coalescence of droplet is feasible to happen because of the mutual effect (or interaction) among droplets in the continuous phase flow. Between these two droplets, also, the coalescence will be successful if the trapped film is presumably drained out completely. Also, several closure models are proposed for the coalescence source terms in the literature [8].

Unlike the other existing PB models, Prince and Blanch [11] model have been developed in a conceptual way based on a multi-stage coalescence mechanism. Furthermore, this model has been applied and verified by many other studies [28, 29]. On the other hand, the equations of Prince and Blanch [11] model have been assembled based on a summation concept to consider different effects (shear, buoyancy, and turbulence) on collision rate. These specifications make this model a multi-factor model, which enable to amount the effect of new factors, like non-hydrodynamic ones to modify a multi-physics model. Therefore, Prince and Blanch [11] model is used in this study to sum electrostatic contributions and modify an EHD PB model. The original model equations and the modifications are presented in the section of "Hydrodynamic Part of PB Model."

The coalescence rates in Equations 4 and 5 explain the possibility of complete collisions between two pairs of droplets, i.e. V and V' . It is determined by multiplying one quantity to the other quantity, which it means the frequency of collision (θ) and the efficiency of coalescence η , as seen in Equation 6.

$$c(V,V')=\theta(V,V') \eta(V,V') \quad (6)$$

Furthermore, the death and birth rates, D^B and B^B , of droplets of volume V because of breakage (in Equation 3) are given by the following equations:

$$B^B(V,t)=\int_{\Omega_V} m b(V')\beta(V|V')n(V',t) dV' \quad (7)$$

$$D^B(V,t)=b(V) n(V,t) \quad (8)$$

where m is number of daughter droplets, $b(V')$ is the frequency of breakup, that is the percentage of particles of volume V' breaking in unit time; in addition, $\beta(V|V')$ is the probability density function (PDF), which it explains the distribution or size range of the daughter droplets.

Closure models for the break-up rate of droplets or bubbles have been proposed in some literature [13]. Because of a very laminar nature of the flow in the electrocoalescers, probability of droplet breakup is very low in this case. Nevertheless, to generalize the PBM, by implementing and evaluating some models, the well-known model developed by Prince and Blanch [11] has been selected to apply as the breakage model in this work. In addition, the related equations of this model are presented in the section of "Hydrodynamic Part of PB Model". The term of rate of breakage composes (1) the frequency of breakage (b), which indicates the percentage of droplets that split in unit time and (2) the probability density function (PDF), which explains the size range of the daughter droplets that have been propagated or come from the splitting mother particle. The term of rate of breakage in Equation 7 is determined by using the following equation:

$$B^B(V,t)=b(V') \beta(V|V') \quad (9)$$

The next sections define how the various functions in the source expression (or terms) are modeled in the present study to couple hydrodynamics and electrostatics in crude oil electrocoalescers.

Hydrodynamic Part of PB Model

The main subject of this study is to model droplet-droplet collision and coalescence either forced by fluid flow or caused by the electrostatic forces in electrocoalescers. The existing PB models only consider hydrodynamic aspects and are not concern about electrostatic issues. As it has been mentioned above, the well-known model developed by Prince and Blanch [11] has been selected to apply as hydrodynamic PB model which is modified in such a way to add the electrostatic effects.

A fundamental model for the bubble coalescence rates and breakage in gas-liquid mixtures have been proposed by Prince and Blanch [11]. Moreover, the coalescence has been modeled by considering collisions due to turbulence, buoyancy, and laminar shear, and by analysis of the coalescence efficiency of collisions. Moreover, the breakup in terms of bubble interactions with turbulent eddies has been analyzed.

Also, it has been suggested by Prince and Blanch [11] that the coalescence of two droplets happen in three steps. First, the two droplets collide, and then a small amount of continuous phase between them is trapped. This liquid is then drained until the liquid film separating the droplets reaches a critical thickness. At this point, film rupture occurs, and this occurrence results in coalescence. With a view to defining whether a considered collision will cause coalescence or not, it is essential that the efficiency of collision be determined. Moreover, two droplets will coalesce if they stay in contact for a period of enough time for the drainage of the trapped liquid film to achieve the critical (or severe) value necessary for a break.

From the hydrodynamic aspect, collisions may happen because of a variety of mechanisms consisting of turbulence, buoyancy, and laminar shear. Droplets may collide by the fluctuating turbulent velocity of the continuous phase. The frequency of collision which comes from turbulent motion (θ_{ij}^T) can be determined as a function of droplet size, velocity, and concentration, by using the following equation [11]:

$$\theta_{ij}^T = n_i n_j S_{ij} (\bar{u}_t^2 + \bar{u}_{tj}^2)^{1/2} \quad (10)$$

where n_i and n_j are the number densities of droplets of radius r_i and r_j respectively; moreover, \bar{u}_t is the average turbulent fluctuating velocity

of the droplet, and S_{ij} which is the collision cross-sectional area of the droplets is defined by using the following equation [11]:

$$S_{ij} = \frac{\pi}{4} (r_i + r_j)^2 \quad (11)$$

Buoyancy-driven collisions which come from the difference in rising velocities of droplets of different size. The buoyant collision rate (θ_{ij}^B) is presented by using the following equation [11]:

$$\theta_{ij}^B = n_i n_j S_{ij} (u_{r_i} + u_{r_j}) \quad (12)$$

where u_r is the rise velocity which can be determined as a function of droplet size. An expression for rising velocity is given by using the following equation [11]:

$$u_{r_i} = [(1.07\sigma) / \rho_o r_i + 1.01g r_i]^{1/2} \quad (13)$$

where σ is the surface tension, and ρ_o is the continuous phase density.

The third contribution to the hydrodynamic collision rate is caused by laminar shear in the continuous phase. Moreover, collisions occur in this situation as a result of the development of a gross circulation pattern in the vessel at high flow rates. The functional form of the collision rate because of laminar shear θ_{ij}^S is given by using the following equation [11]:

$$\theta_{ij}^S = n_i n_j \frac{4}{3} (r_i + r_j)^3 \gamma \quad (14)$$

where γ is the average shear rate.

With the intention of specifying what fraction of droplet collisions leads to coalescence phenomena, it is indispensable that collision efficiency be determined. The efficiency is a function of the contact time between droplets and the necessary time for droplets to coalesce. A statement for the coalescence efficiency An expression for the coalescence efficiency (η_{ij}) is given by using the following equation [10]:

$$\eta_{ij} = \exp(-t_{ij} / \tau_{ij}) \quad (15)$$

where t_{ij} is the time required for the coalescence

of droplets of radius r_i and r_j , while τ_{ij} is the contact time for the two droplets. Moreover, the coalescence time is given by using the following equation [11]:

$$\tau_{ij} = \left(\frac{r_{ij}^3 \rho_o}{16\sigma} \right)^{1/2} \ln \frac{h_o}{h_f} \quad (16)$$

where h_o is the initial film thickness and h_f is the critical film thickness where rupture occurs; in addition, r_{ij} is the equivalent radius of droplets of radius r_i and r_j which is given by using the following equation [11]:

$$r_{ij} = \frac{1}{2} \left(\frac{1}{r_i} + \frac{1}{r_j} \right)^{-1} \quad (17)$$

The time that droplets remain in contact is dependent on the droplets size and the collision intensity. High levels of turbulence increase the probability that an eddy will separate droplets, while large droplet size provides larger contact areas. An estimate of the contact time is given by using the following equation [11]:

$$\tau_{ij} = \frac{r_{ij}^{2/3}}{\epsilon^{1/3}} \quad (18)$$

Values for the coalescence and contact time may be substituted into Equation 15 to determine the efficiency. According to Equation 6, the overall hydrodynamic coalescence rate of droplets of radius r_i and r_j is defined by multiplying the total hydrodynamic collision frequency by the coalescence efficiency (Equation 19).

$$c(r_i, r_j) = \frac{1}{2} (\theta_{ij}^r + \theta_{ij}^s + \theta_{ij}^s) \exp(-t_{ij} / \tau_{ij}) \quad (19)$$

The factor 1/2 is included to avoid counting coalescence events between droplet pairs twice. A model for breakage rates of liquid droplets based on the fraction of droplets undergoing breakup and the time required for breakage to occur has been developed by Coulaloglou and Tavlarides [10]. In addition, it is assumed that the breakup rate is a function of the dispersed phase density.

Subsequently, a model for breakup based on the interaction of droplets with turbulent eddies has been developed by Prince and Blanch [11]. Moreover, the turbulent collision rate of droplets with eddies of the appropriate size has been considered by them, and the rate is given by using the following equation [11]:

$$\theta_{ie} = n_i n_e S_{ie} (\bar{u}_{ti}^2 + \bar{u}_{te}^2)^{1/2} \quad (20)$$

which is analogous to Equation 10, while some parameters are replaced with the eddy diameter, concentration, and velocity. To apply this equation, the number of eddies of a specified size must be identified. It is assumed by Prince and Blanch [11] that the isotropic turbulence is governed, and in the inertial subrange, the eddy size of interest lies. Furthermore, only a determined amount of collisions between droplet and eddy causes in droplet breakup probably. Moreover, the criterion or limitation for breakage depends on the kinetic energy of the eddy to the surface tension forces of the droplet, which is expressed in terms of the critical Weber number by Prince and Blanch [11]. Accordingly, they have defined collision efficiency for breakup has been defined by them by using the following equation [11]:

$$\beta = \exp(-u_{ci}^2 / u_{te}^2) \quad (21)$$

where u_{ci} is the critical eddy velocity for the breakup of a droplet of radius r_i , and u_{te} is the turbulent velocity of an eddy of radius r_e .

Finally, the total breakup rate has been presented by multiplying Equation 20 by Equation 21 as seen in Equation 22 [11]:

$$\sum_i \sum_e \theta_{ie} \exp(-u_{ci}^2 / u_{te}^2) \quad (22)$$

Electrostatic Part of PB Model

The hydrodynamic part of PBM is presented in the section of Hydrodynamic Part of PB Model based on the well-known model developed by Prince and Blanch [11]. In this section, the electrostatic issues of the PBM are described based on an electrostatic model developed by Atten [26].

The basic process in the electrocoalescences is the interaction between neighboring droplets because of the mutual attraction of dipoles induced by the electric field, and there by electrocoalescence is resulted. Furthermore, an infrequent process relates to the feasible elongation of the droplets under the effect of a high enough electric field, $E_c = 0.64(\sigma / 2\epsilon)^{1/2}$, and the subsequent breakage into the smaller droplets [14-16].

The equation which governs the progress of a poly-dispersed emulsion has been derived by Atten [26]. The presented expressions have been obtained from the coalescence rate when only the electric field induces interaction between droplets. In the case of two droplets of radius r_i and r_j in a uniform field E_0 , the interaction force can be calculated as the dipole-dipole force. If there is a skew angle between the direction of the electrical field and the axis of the drops, the dipole-dipole force separates into two components. The radial component force, F_r , moves the drops towards each other and the tangential component force, F_θ , induces a torque on the drops to align them with the field direction. The force components have been presented by using the following equations:

$$F_r = -12\pi\epsilon_0 r_j^3 E_0^2 (r_i^3 d^{-4}) (3\cos^2 \Psi - 1) \quad (23)$$

$$F_\theta = -12\pi\epsilon_0 r_j^3 E_0^2 (r_i^3 d^{-4}) \sin(2\Psi) \quad (24)$$

where ϵ_0 is the permittivity of the oil phase.

A rough estimation method to determine the

characteristic time for the evolution of the mean droplet size of an emulsion exposed to an electric field has been applied by Atten [26]. A basic stage for transforming N drops of radius to $N/2$ drops of radius $(2^{1/3}r)$ has been considered by Atten. The characteristic time has been identified from resolution of the equation which governs the relative movement of two adjacent droplets. Through a simplified assumption by using Stokes' formula for the viscose friction and the dipolar expression:

$$4\pi v_o r \frac{ds}{dt} = 2F = -48\pi\epsilon_0 E_0^2 (r^6 d^{-4}) \quad (25)$$

where v_o is the dynamic viscosity of the oil phase. Consequently, the characteristic time for doubling of droplet radius has been calculated by using the following equation [26]:

$$t_{ij}^E = (8/5) \left(v_o / \epsilon_0 E_0^2 \right) \left[(\pi / 6\Phi)^{5/3} - 1 \right] \quad (26)$$

where Φ is the volume fraction of the water phase.

A mono-dispersed emulsion to obtain some orders of magnitude regarding the time evolution without numerical integration has been assumed by Atten [26]. By these assumptions, the rate of collision for drops of radius r_i and r_j has been obtained:

$$\theta_1^E = n_i n_j (16\epsilon / 3\sqrt{3}) (\epsilon_0 E_0^2 / v_o) r_i^2 r_j^2 / (r_i + r_j) \quad (27)$$

With Equation 27, the characteristic time of first decay of a mono-dispersed emulsion can be estimated by using the following equation [29]:

$$t_{ij}^E = (\sqrt{3} / 4\Phi) (v_o / \epsilon_0 E_0^2) \quad (28)$$

For small distances between drops, the viscous force which slows the relative motion of drops cannot be obeyed the Stokes' law usually. Braking the motion originates from the thinning of the oil phase which lies among the drops. For these conditions, the collision rate has been derived by using the following equation [26]:

$$\theta_2^E = n_i n_j (2\pi / 9) (r_i + r_j)^3 (r_i / r_j + r_j / r_i) (\epsilon_0 E_0^2 / v_o) \quad (29)$$

It has been suggested by Atten [26] that however Equations 27 and 29 give a good estimate for the kinetics of electrocoalescence, some explicit and implicit estimation such as assuming homogeneous and isotropic spatial distribution of drops, as well as the simplified hydrodynamic aspect, limit this model to generalize directly for practical applications. Therefore, a modification can be achieved by combination of electrostatic PB model of Atten [26] with hydrodynamic PB model of Prince and Blanch [11].

Electrohydrodynamic PB Model

To develop an electrohydrodynamic population balance model, the governing equations shall be modified to consider the effect of the electric field both in the collision frequency and the coalescence efficiency. The model shall be developed so that the other mechanisms of the collision caused by the hydrodynamic, e.g. turbulence, buoyancy, and laminar shear can be simultaneously contributed in the collision frequency term beside the electric field. Furthermore, the model shall be able to decide the governing coalescence mechanism to calculate the coalescence efficiency. For this purpose, the hydrodynamic PB model by Prince and Blanch [11] is combined with the electrostatic PB model by Atten [26].

As it is mentioned in the section of "Electrostatic Part of PB Model", two expressions for electrostatic collision frequency have been presented by Atten: (1) θ_1^E (Equation 27) from a simplified assumption by using the formula of Stokes for the viscose friction in balance with the dipolar force, (2) and θ_2^E (Equation 29) by using the film thinning force which provides more accurate results [29]. The latter expression will be applied in this study as the

electrostatic collision frequency.

Moreover, collisions from different mechanisms are accumulative as seen in Equation 19. Therefore, the collision frequency in the modified EHD PB model is calculated using the following equation:

$$\theta(r_i, r_j) = \frac{1}{2} (\theta_{ij}^T + \theta_{ij}^B + \theta_{ij}^S + \theta_{ij}^E) \quad (30)$$

where θ_{ij}^E is the electrostatic collision frequency from Equation 29.

On the other hand, as it is mentioned in the section of "Hydrodynamic Part of PB Model," when an emulsion is exposed to a shear, the droplets which have been dispersed have a relative motion, and several collisions happen. Also, the respective collision frequency has been derived for a Couette flow as presented in Equation 14. This relation infers that every collision causes that coalescence of the colliding droplets occurs [15]. Without electric field, this is not true because there is a quite long procrastination between droplets contact and coalescence. Also, the respective characteristic time for laminar shear has been derived by Atten as seen in Equation 31:

$$t_{ij}^S = \pi / (16 \phi \gamma) \quad (31)$$

According to Equation 15, the efficiency of coalescence will be a function of the contact time between droplets and the time needed for droplets to coalesce.

In order to calculate the efficiency in the modified EHD PB model, the governing coalescence mechanism and respective characteristic time shall be determined. The characteristic time for the hydrodynamic coalescence mechanism is given by Equation 31, and the characteristic time for the electrostatic coalescence mechanism is given by Equation 28. Physically, the coalescence is completed in the quicker time of these two

mechanisms. Therefore, the characteristic time in the modified EHD PB model is determined as the minimum between these two values by using the following equation (Equation 32), and it is substituted in Equation 15.

$$t_{ij} = \min(t_{ij}^E, t_{ij}^S) \quad (32)$$

Accordingly, the characteristic time in the spatial position of the droplets is calculated at each time step. Consequently, the overall coalescence rate of droplets will be computed by the all of frequency of collision (Equation 30) multiplied by the efficiency of coalescence (Equation 15) as seen in Equation 33:

$$c(r_i, r_j) = \theta(r_i, r_j) \exp(-t_{ij} / \tau_{ij}) \quad (33)$$

Numerical Simulation Procedure

Different numerical methods have been proposed as solutions to PB equations like the method of moments, weighted residuals, discrete formulations, and Monte Carlo simulation methods. Some of these methods have been reviewed and categorized before [27]. In discretized methods, all particles of the same size are called a class. These methods divide the size range into small cells and then apply a balance equation for each cell. The continuous population balances are reduced to a set of ordinary differential equations. An efficient technique to solve PBEs is called fixed-pivot. Moreover, it has been proposed by Kumar and Ramkrishna [28], and then its accuracy and convergence have been improved and renamed to cell average technique [29].

In this study, the equations of population balance are solved by the discrete or class method (CM) [30]. In the CM, the particle population is discretized into a finite number of intervals which are used to track the population density directly. To achieve a complete CFD-PBM coupling, the distribution of the droplet size is separated into some of discrete

size classes, while it is presumed that droplets of all sizes share a usual velocity field related to the water phase.

The PBE (Equation 2) is rewritten in terms of volume fraction of droplet class i for the discrete method as seen in Equation 34:

$$\frac{\partial}{\partial t} [\rho_w \Phi_i] + \nabla \cdot [\rho_w u_i \Phi_i] = \rho_w V_i S \quad (34)$$

where Φ_i is the volume fraction of droplet class i , and it is defined by using the following equation:

$$\Phi_i = N_i V_i = V_i \int_{V_i}^{V_{i+1}} n(V, t) dV \quad (35)$$

Furthermore, a fraction of Φ called f_i is applied as a solution variable. This fraction is defined by using the following equation:

$$f_i = \frac{\Phi_i}{\Phi} \quad (36)$$

where Φ is the total volume fraction of the water phase.

The particle birth and death rates of Equations 4 and 5 and Equation 7 and 8 are rewritten in terms of droplet class i or the discrete method as follows:

$$B_i^C = \sum_{k=1}^N \sum_{j=1}^N c_{kj} N_k N_j \delta_{kj} \omega_{kj} \quad (37)$$

$$D_i^C = \sum_{j=1}^N c_{ij} N_i N_j \quad (38)$$

$$B_i^B = \sum_{j=i+1}^N b_j N_j \beta(V_i | V_j) \quad (39)$$

$$D_i^B = b_i N_i \quad (40)$$

where $c_{ij} = c(V_i, V_j)$ and the parameters ω and δ preserve the reassignment to the classes i , the resulting coalesced droplet lies among the classes.

$$\omega_{kj} = \begin{cases} 1 & \text{for } V_i < V_c < V_{i+1}, \quad i \leq N-1 \\ 0 & \text{elsewhere} \end{cases} \quad (41)$$

V_c is the droplet volume resulting from the coalescence of droplets k and j , and is defined by using the following equation:

$$V_c = [\delta_{kj} V_i + (1 + \delta_{kj}) V_{i+1}] \quad (42)$$

where

$$\delta_{kj} = \frac{V_c - V_{i+1}}{V_i - V_{i+1}} \quad (43)$$

Furthermore, for coupling PBM with the fluid dynamics, the Sauter mean diameter of the droplets can be calculated by using the following equation:

$$d_{32} = \frac{\sum N_i d_i^3}{\sum N_i d_i^2} \quad (44)$$

Fluid Flow Governing Equation

The Eulerian-Eulerian method is most constructive for the present liquid-liquid flow, and thereby it was used. For each phase, the conservation equations of continuity and momentum are solved. The continuity equation is, ($i=o,w$):

$$\frac{\partial}{\partial t}(\alpha_i \rho_i) + \nabla \cdot (\alpha_i \rho_i \vec{u}_i) = 0 \quad (45)$$

where α is the volume fraction, which indicates the space filled up by each phase, ρ is the phase density, and u is the phase velocity. The conservation equation of momentum is:

$$\frac{\partial}{\partial t}(\alpha_i \rho_i \vec{u}_i) + \nabla \cdot (\alpha_i \rho_i \vec{u}_i \vec{u}_i) - \nabla \cdot \tau_i = -\alpha_i \nabla p + \alpha_i \rho_i \vec{g} + F_k \quad (46)$$

where τ is the stress-strain tensor, p is the pressure which is common between all phases, g is the gravitational acceleration, and F_k reveals the interfacial forces. The inter-phase interaction term F_k consists of different momentum exchange mechanisms. For example, drag force is taken into account, and the model of Schiller and Naumann is applied. Moreover, for the volume fractions, the limitations must also be adjusted:

$$\alpha_o + \alpha_w = 1 \quad (47)$$

The two-way coupling between CFD and PBM is undertaken in the model. This means that the result of the flow simulation is applied for obtaining the

drop size distributions. Then again the effect of the drops is assessed on the flow field. The drop size distribution is separated into several discrete size classes, while it is presumed that droplets of all sizes have a shared velocity field, where the Sauter mean diameter has an effect on the drag force. The Sauter mean diameter of the droplets is obtained in each time step (Equation 44) by PBM and retrieved to CFD to calculate the drag force. All drop size classes have the same velocity field as the water phase, therefore Navier–Stokes equations is solved for all drop classes. The required parameters like turbulent energy dissipation in the source terms of PBM for coalescence and breakage is retrieved for each computational cell from the CFD. The velocity and pressure fields of the continuous oil phase and the dispersed water phase are obtained from the Navier–Stokes equations using the Eulerian approach.

Regarding to the number of classes in the CM method, the precision is improved at the cost of rised CPU time. As the calculating time is restrictive, a balance has to be established between CPU time and the selected number of classes [17]. Several simulations have been executed using the number of different classes of 8, 10, 15, and 20. Then the results have been compared with each other to decide about the reasonable case. Finally, the results have demonstrated that using 10 classes of droplet sizes is acceptable for our case.

Problem Discription

The modified model is implemented for the same geometry and the same applied conditions in a pilot electrocoalescer to evaluate the validity of the EHD model in comparison with the experimental data. The pilot data have been attained in a continuous-

flow horizontal drum as the electrocoalescer. Also, owing to the axial symmetry of the system, a circular cross-section of the vessel is considered as the geometry of the problem (Figure 1). The diameter of the used electrocoalescer as shown in Figure 1 is 38 cm.

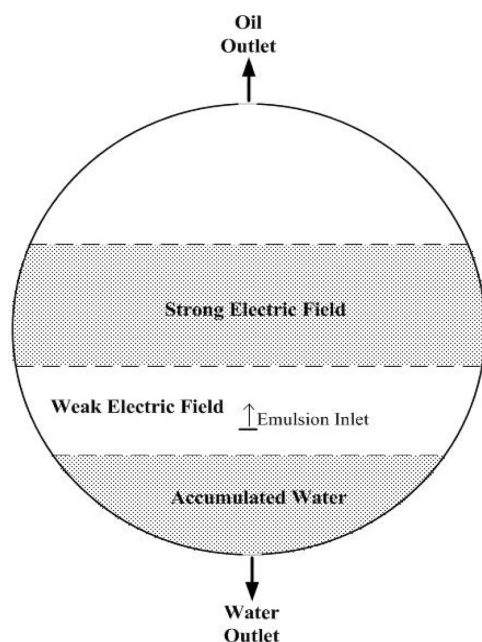


Figure 1: The geometry of the electrocoalescer.

The applied operating conditions in the pilot test are given in Table 1.

Table 1: Applied operating conditions in the pilot test.

Temperature (°C)	Pressure (Barg)	Inlet Water Content (vol. %)	Electric Field Intensity (kV/cm)
70	7	5.5	1.4

The physical properties of the used fluids at the operating conditions have been presented in Table 2.

Table 2: Physical properties of the phases.

Phase	Density (kg/m ³)	Viscosity (cP)
Oil	818	4.4
Water	985	0.45

Moreover, the drop size distribution of water phase in the inlet emulsion is given in Table 3.

Table 3: Droplet size distribution of the inlet emulsion.

Size Group Number	Droplet Size (μm)	Fraction (vol. %)
S ₁	2	2
S ₂	4	2
S ₃	8	5
S ₄	16	10
S ₅	32	20
S ₆	64	30
S ₇	128	15
S ₈	256	10
S ₉	512	5
S ₁₀	1024	1

The inlet velocity of the emulsion is 6 mm/s regarding the inlet flow rate of the pilot electrocoalescer.

The equations of the model construct a transient problem; and therefore, a proper unsteady state solution approach is applied. The time-dependent PB-CFD problem has been solved using the aforementioned equations of the model. The variations in the distribution of different size groups of the droplets have been inspected through the computational domain in progress run-time to study the droplet size growth and water separation until the steady state is reached.

RESULTS AND DISCUSSION

Numerical Results

The EHD model has been executed for the above-mentioned problem in the section of “Problem Description.” Also, results are presented at several times: $t = 30, 60, 90, 120, 150$, and 180 min in Figures 2 to 7. The predicted distributions of water phase are presented in Figure 2 from the CFD-PB simulation.

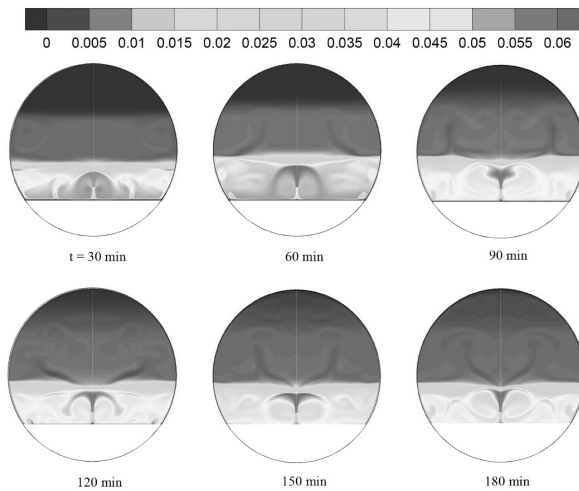


Figure 2: Fraction of dispersed water phase in the oil phase.

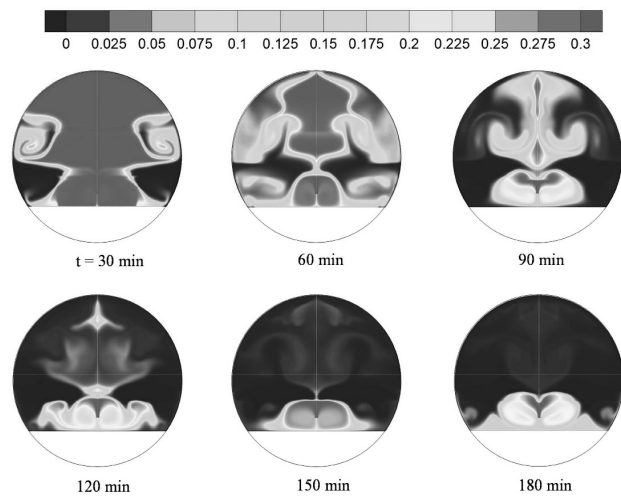


Figure 5: Fraction of S_6 size (64 μm) droplets in dispersed water content of the oil phase.

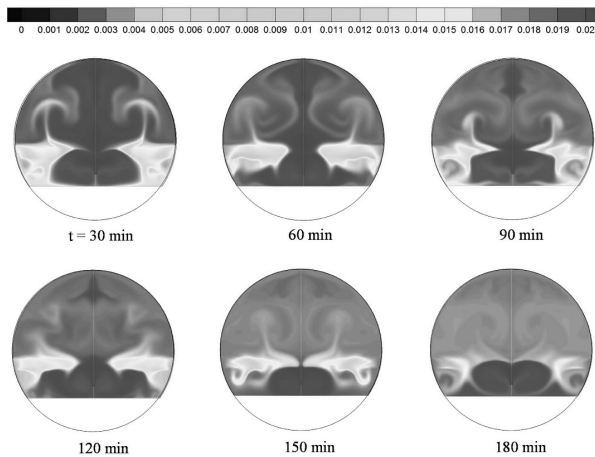


Figure 3: Fraction of S_2 size (4 μm) droplets in dispersed water content of the oil phase.

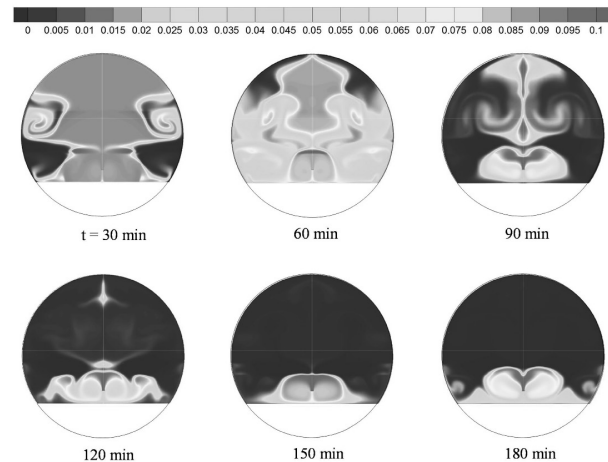


Figure 6: Fraction of S_8 size (256 μm) droplets in dispersed water content of the oil phase.

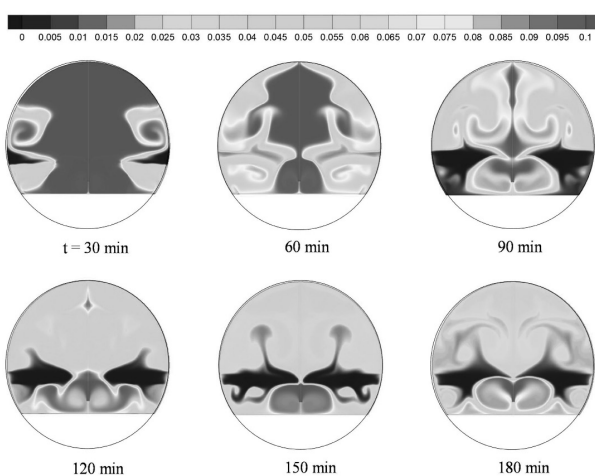


Figure 4: Fraction of S_4 size (16 μm) droplets in dispersed water content of the oil phase.

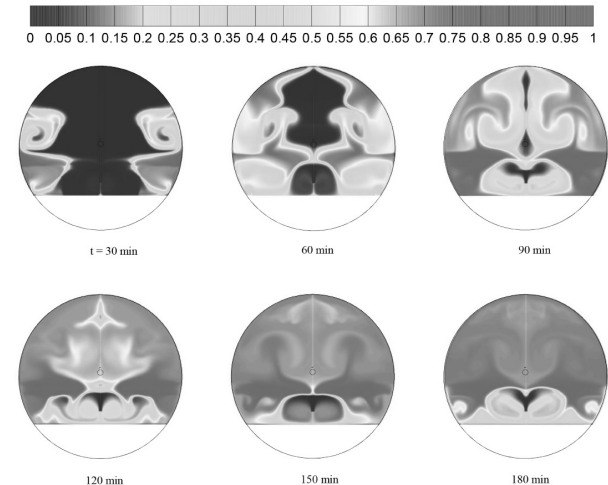


Figure 7: Fraction of S_{10} size (1024 μm) droplets in dispersed water content of the oil phase.

According to Table 1, the water content of the inlet emulsion is 5.5 vol. %. At the time of 30 min, a thick region of emulsion with the water content of 4-5% has been formed above the accumulated water phase. A narrow area with the water content of 2-3% is seen above the said region. Next, a thick region of the dilute emulsion has been accumulated with the water content of about 1%. This region corresponds to the strong electric field area that seems effective in separating the water droplets. Finally, the upper region with very low water content represents the oil phase passed through the strong electric field.

The aforementioned arrangement is relatively the same at the time of 60 and 90 min, but the region of the dilute emulsion is developing towards the upside of the vessel. At $t = 120$ min, no significant changes were seen in the bottom half of the vessel. However, in the upper half of the vessel, the height of the dilute emulsion region has been increased upward, and instead, the thickness of the upper region has been decreased. The above-mentioned trend is continued at $t = 150$ min and $t = 180$ min. Since the profiles are not changed even up to $t = 240$ min, the solution is considered as the steady-state condition at the time of 180 min.

In Figure 3, the profiles of S_2 size ($4\text{ }\mu\text{m}$) droplets from the CFD-PB simulation at several times are displayed.

According to Table 3, the fraction of S_2 size in water phase of the inlet emulsion is 2 vol%. At $t = 30$ min, the fraction of S_2 size remained the same as inlet emulsion in the upper half of the vessel. However, the fraction of S_2 size was decreased to about 1-1.5% through electrocoalescence in some area of the bottom half. The mentioned condition

is relatively the same at $t = 60$ min, but the area of the decreased fraction in the bottom half is limited. However, at this time, the fraction had started to decrease through electrocoalescence in the upper half of the vessel to about 1.7%. The aforesaid trend is continued at $t = 90$ to 180 min. A uniform profile was shaped at $t = 180$ min, as the fraction in the upper half of the vessel was reduced to about 1.6%, and the fraction in the bottom half was remained about 2%. Ultimately, similar results are seen in the profiles of S_1 size ($2\text{ }\mu\text{m}$) droplets from the CFD-PB simulation.

The profiles of S_4 size ($16\text{ }\mu\text{m}$) droplets are presented in Figure 4 from the simulation results. As presented in Table 3, the fraction of S_4 size droplets in water phase of the inlet emulsion is 10%. At $t = 30$ min, the fraction of S_4 size remained the same as inlet emulsion in major part of the vessel. However, the fraction of S_4 size was decreased to about 5% through electrocoalescence in some area near the wall. The area of the decreased fraction was extended at $t = 60$ min, except for the central region. However, at $t = 90$ min, more progress of the electrocoalescence reduced the fraction of S_4 size, even in the aforementioned central area of the vessel. The decreasing trend of the fraction was continued, while a constant profile was formed at $t = 180$ min. Approximate similar results were observed for the profiles of S_3 size ($8\text{ }\mu\text{m}$) and S_5 size ($32\text{ }\mu\text{m}$) droplets from the simulation.

Figure 5 displays the profiles of S_6 size ($64\text{ }\mu\text{m}$) droplets resulted from the CFD-PB simulation. As given in Table 3, the fraction of S_6 size droplets in water phase of the inlet emulsion is 30%. At $t = 30$ min, the fraction of S_6 size remained equal to the inlet content in major part of the electrocoalescer.

However, it was decreased to about 0-10% in some area near the wall. Moreover, consumption of S_6 size droplets through electrocoalescence was extended at $t=60$ min, aside from the central section. At $t=90$ min, more development of the electrocoalescence reduced the fraction of S_6 size droplets in a major area of the vessel to about zero, significantly in the central section. At the next times, this reduction was continued to form a uniform profile with about zero S_6 size droplets, except for a limited region around the inlet nozzle. Moreover, approximate similar consequences resulted from the simulation for the profiles of S_7 size ($128\ \mu\text{m}$) droplets.

The profiles of S_8 size ($256\ \mu\text{m}$) droplets are shown in Figure 6 from the simulation results.

The fraction of S_8 size droplets in water phase of the inlet flow is 10%, as presented in Table 3. At $t=30$ min, the fraction of S_8 size was decreased from 10% to about 9% in the major part of the vessel. However, the fraction was more reduced to about 0-4% in some area near the wall. At $t=60$ min, consumption of S_8 size droplets was accelerated in the entire of the vessel. At $t=90$ min, more progress of the electrocoalescence reduced the fraction in a major area of the vessel to about zero, except for a narrow region of the central section. At the next times, S_8 size droplets were only observed in a limited region around the inlet nozzle, which the droplets were consumed gradually by electrocoalescence. Approximate similar results were observed for the profiles of S_9 size ($512\ \mu\text{m}$) droplets from the simulation.

Figure 7 shows the profiles of S_{10} size ($1024\ \mu\text{m}$) droplets resulted from the CFD-PB simulation. As presented in Table 3, the fraction of S_{10} size droplets

in water phase of the inlet emulsion is 1%. At $t=30$ min, the fraction of S_{10} size remained equal to the inlet content in major part of the electrocoalescer. Though, it was increased to about 50-90% in some area near the wall. Formation of S_{10} size droplets through the electrocoalescence of smaller droplets was extended at $t=60$ min, except for the central section. At $t=90$ min, more development of the electrocoalescence increased the fraction of S_{10} size droplets in the entire of the vessel significantly. Increasing the fraction was accelerated at $t=120$ min and continued to form a uniform profile with about 70-90% S_{10} size droplets. However, in a limited region around the inlet nozzle, the fraction was less, in which S_{10} size droplets were generated progressively from electrocoalescence of the smaller droplets.

Model Validation

As seen in Figure 2, the CFD-PB model predicts the water content in the outlet oil stream equal to 0.33% for the steady state conditions at $t=180$ min. On the other hand, the result of the pilot test for the same operating condition reveals this value as 0.29%. Therefore, the predicted value of the developed PB model is validated in comparison with the experimental data. The difference between the results of the model with the actual value may be explained regarding some assumptions in the model and some probable inaccuracy of the experimental sampling and measurements.

RESULTS AND DISCUSSION

The modified electrohydrodynamic PB model was implemented in a CFD simulation applying the same operating condition of the completed pilot test. The resulted profiles of water distribution in

Figure 2 demonstrate the effect of electric field on water-oil separation performance. In the weak electric field zone between water accumulation zone and strong electric field zone (see Figure 1), the inlet w/o emulsion enters with 5.5 vol. % water content and a specified droplet size distribution (Table 3).

According to Equations 10, 12, 14, and 29, high water content, and thereby high number density of droplet size groups, enhances the probability of drops collision rates in this zone. The profile of water content shows an increasing trend from the top to the water-oil interface. On the opposite direction, the upward vertical flow of the emulsion with less water content enters the strong electric field zone. High electric field strength increases the probability of drops electrical collision rate (Equation 29) and coalescence efficiency (Equation 28). The profile of water content shows a decreasing trend from the strong electric field to the top of the vessel. Furthermore, by comparing numerical results and the experimental data, it is demonstrated that the effect of the electric field in the separation of water from the oil is predicted successfully by the EHD PB model.

The resulted profiles of some droplet size groups have been presented in Figures 3 to 7. The distributions of S_2 (Figure 3) and similarly, for S_1 size groups illustrate that the applied electric field intensity has a limited effect on these very small size (2-4 μm) droplets. Moreover, this result can be also explained regarding Equations 11, 14, and 29, as the collision rates and coalescence efficiency decrease significantly for the smaller droplets. The comparison between Figure 3 with Figures 4, 5, and 6 demonstrates that the removal of very small water droplets (S_1 - S_2) through electrocoalescence

is negligible in comparison with the bigger droplets. Furthermore, the industrial applications of crude oil electrocoalescers confirm that the separation of very small size droplets is the most important challenge of this process.

The comparison of the distributions of S_4 (Figure 4), S_6 (Figure 5), and S_8 (Figure 6) size groups with those of S_{10} (Figure 7) size group illustrates a direct relevance among the profiles at the same times. For example, at $t = 30$ min, the observed patterns for S_4 , S_6 and S_8 in Figures 4 to 6 showed specified decreased fraction regions which were coincided with the increased fraction areas of S_{10} size group in Figure 7. Moreover, this analogy is seen for the other displayed times similarly. These analogies show that reduction in the fraction of the small droplets finally results in an increase in the fraction of the biggest droplet. Therefore, the consequences confirm that electrocoalescence is successful for separating the water through the growth of the droplet size.

On the other hand, by comparing the results in Figures 3 to 6, some differences are observed among the variations of different drop sizes. For the very small droplets (S_1 - S_2), the reduction of the fraction is negligible (Figure 3), while for the medium droplets (S_3 - S_5), the decrease is more considerable, but these are not completely removed by electrocoalescence (Figure 4), whereas for the larger droplets (S_6 - S_9), the reduction is so significant (Figure 5 and 6) as these droplets have been eliminated completely through electrocoalescence.

The resulted profiles of S_{10} size droplets (Figure 7) illustrate that the fraction of this size group shows the largest variations among the whole size groups as its near-zero fraction in the inlet flow is increased significantly in the major parts of the

vessel to about 70-90%. Additionally, in the region below the strong electric field, where the formed big droplets settle down, the most fractions of S_{10} size droplets become about 98%.

CONCLUSIONS

A special electrohydrodynamic PB model has been developed, which the developed model considers both hydrodynamic and electrostatic aspects in w/o emulsions. For this purpose, the hydrodynamic PB model developed by Prince and Blanch [11] was combined with the electrostatic PB model developed by Atten [26]. The overall coalescence rate of droplets is defined by the all of frequency of collision multiplied by the coalescence efficiency. Also, the model was developed so that the mechanisms of collision from the hydrodynamic sources, i.e. buoyancy, turbulence, and laminar shear are contributed in the collision frequency term beside the electric field simultaneously. Furthermore, the coalescence efficiency is calculated regarding to the govern mechanism. The modified model was implemented in a combined CFD-PBM simulation applying the same operating condition of the completed pilot test. The resulted profiles of water phase distribution demonstrated the effect of electric field on water-oil separation performance. In the zones with high water content and consequently high number densities of droplets, more progress in electrocoalescence was observed because of the enhanced probability of the collision rate. The profile of water content showed an increasing trend from the top to the water-oil interface. The probability of the electrical collision rate and coalescence efficiency was enhanced in the upward vertical flow of the emulsion inside the

strong electric field zone. Therefore, a decreasing trend was observed in water content from the strong electric field to the top. The predicted water content in the outlet oil (0.33%) was in a good agreement with the experimental data (0.29%). Moreover, the profiles of different droplet sizes illustrated that removal of the very small droplets ($<4 \mu\text{m}$) was minor, while reduction of the medium sizes ($8\text{--}32 \mu\text{m}$) was more considerable but not complete whereas elimination of the larger droplets ($>64 \mu\text{m}$) was complete. On the other hand, the largest droplet size ($1024 \mu\text{m}$) displayed a significant population growth, normally to about 90%. Therefore, the modified model assists us in predicting the internal situations of the phases and the size distribution of the droplets in crude oil electrocoalescers. Finally, the resulted data from CFD-PBM simulation are useful for designing the electric field and improving the hydraulic of the internal fluid flow.

NOMENCLATURES

b	: breakage frequency, s^{-1}
B	: birth rates, $\text{m}^{-3} \text{s}^{-1}$
c	: coalescence rate, s^{-1}
D	: death rates, $\text{m}^{-3} \text{s}^{-1}$
d	: distance between drop centers, m
E_0	: electric field intensity, V m^{-1}
F	: force, N
g	: gravity acceleration, m s^{-2}
n	: number density, m^{-3}
p	: the net rate of production of droplets, $\text{m}^{-3} \text{s}^{-1}$
r	: drop radius, m
s	: collision cross-sectional area, m^2
t	: time, s
u	: velocity, m s^{-1}
v	: droplet volume, m^3

Greeks Letters

β	: probability density function
γ	: average shear rate, s^{-1}
ϵ	: energy dissipation rate per mass, $m^2 s^{-2}$
ε	: relative permittivity
η	: coalescence efficiency
θ	: collision frequency, s^{-1}
ν	: dynamic viscosity, Pa. s
ρ	: Density, $kg m^{-3}$
σ	: interfacial tension, $N m^{-1}$
τ	: contact time, s
Φ	: volume fraction
ψ	: electric field skew angle, degree

Subscripts / Superscripts

B	: buoyant
C	: coalescence
E	: electrical
o	: oil
r	: radial direction
S	: laminar shear
T	: turbulent
θ	: tangential direction
w	: water

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